

Consequences of a dressed quark-gluon vertex in heavy-light mesons

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in collaboration with
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FWF Der Wissenschaftsfonds.

P25121-N27

Jefferson Lab, May 21, 2014,

Motivation

- QCD and the structure of hadrons
 - ★ the role of the quark masses in hadron spectrum and structure
- Heavy-light mesons
 - ★ Dynamics of the heavy-light meson controlled by dynamics of the heavy quark
 - ★ For mesons with $m_Q \sim m_M$ and $\frac{m_{\bar{q}}}{m_Q} \sim 0$: some conditions are provided by “heavy-quark symmetry”

Neubert Phys. Rept. 245 (1994)

- Heavy-quark symmetry as a guidance
 - ★ Use constraints of heavy-quark symmetry to check the correct behaviour of the mass
 - ★ Similar procedure within point-form Hamiltonian dynamics using a constituent quark model: fully relativistic treatment of light quark is necessary

Gomez-Rocha, Schweiger PRD 86 (2012)

- Limitations of BSE/DSE description of hadrons when using the rainbow-ladder truncation approach
 - ★ In particular: difficulties to describe heavy-light mesons

The main goal of this work...

Based on previous studies of consequences of a dressed quark-gluon vertex in $q\bar{q}$ mesons...

Bender, Detmold, Thomas, Roberts, PRC **65** (2002)
Bhagwat, Höll, Krassnigg, Roberts, Tandy PRC **70** (2004)

Use the same interaction model:

Munczek and Nemirovsky PRD **28**, 181 (1983)

- sufficiently simple \rightarrow allow to compute an arbitrary number of iterations
- sufficiently realistic \rightarrow important features in common with QCD: e.g. confinement, dynamical chiral symmetry breaking

... attempt to generalize to $q\bar{Q}$ mesons

... ask:

Which kind of consequences do corrections to the bare vertex in heavy-light mesons have?

or...

Which kind of heavy-light physics are we “truncating” using RL approximation?

Dyson-Schwinger/Bethe-Salpeter approach

The interaction model

DSE:

$$S^{-1}(p) = i\not{p} + \mathbf{1}m_q + \int \frac{d^4q}{(2\pi)^4} g^2 D_{\mu\nu}(p-q) \frac{\lambda^a}{2} \gamma_\mu S(q) \Gamma_\nu^a(q;p)$$



Interaction model : $\mathcal{D}_{\mu\nu}(k) := \frac{4}{3} \delta_{\mu\nu} (2\pi)^4 \mathcal{G}^2 \delta^4(k)$

Munczek and Nemirovsky PRD **28**, 181 (1983)

DSE becomes an algebraic equation

⇒ Solvable at every truncation order: $n = 0, 1, 2, 3, \dots$

⇒ Analytical solution in RL

⇒ Solvable in the limit $n \rightarrow \infty$ (fully dressed)

$$\Gamma_\mu^C(q;p) = \text{[Diagrammatic expansion of the vertex function } \Gamma_\mu^C(q;p) \text{ as a sum of tree-level and higher-order diagrams.]}$$

The diagrammatic expansion shows the vertex function $\Gamma_\mu^C(q;p)$ as a sum of terms. The first term is a tree-level vertex (a line with a circle). The second term is a tree-level vertex with a gluon loop. The third term is a tree-level vertex with a gluon loop and a ghost loop. The fourth term is a tree-level vertex with a gluon loop, a ghost loop, and a ghost loop. The expansion continues with higher-order terms.

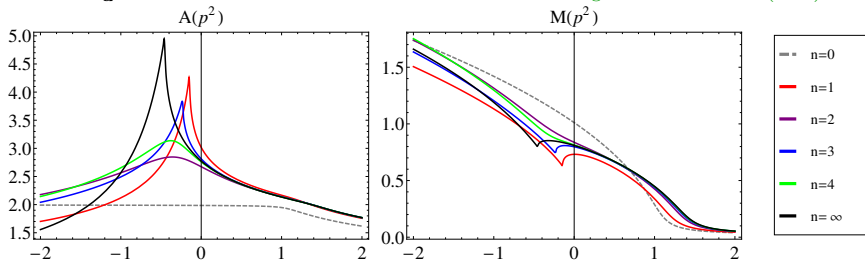
figure adapted from Bender et al. PRC **65** 2002

$$\Gamma_{\mu, n+1}^C(p) = -C \gamma_\nu S(p) \Gamma_{\mu, n}^C(p) S(p) \gamma_\nu$$

Solutions to the DSE for n loops

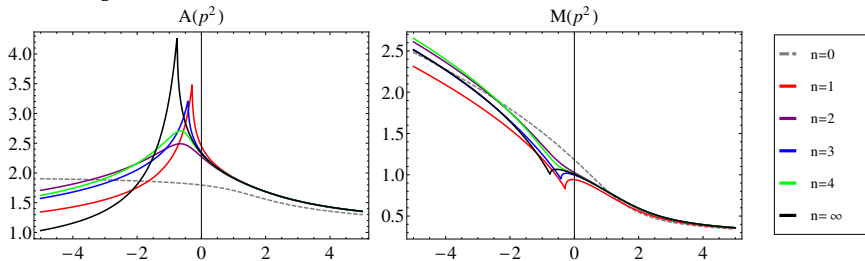
for $m_u = 10$ MeV:

Bhagwat et al. PRC 70 (2004)



for $m_s = 0.166$ GeV:

$C = 0.51$

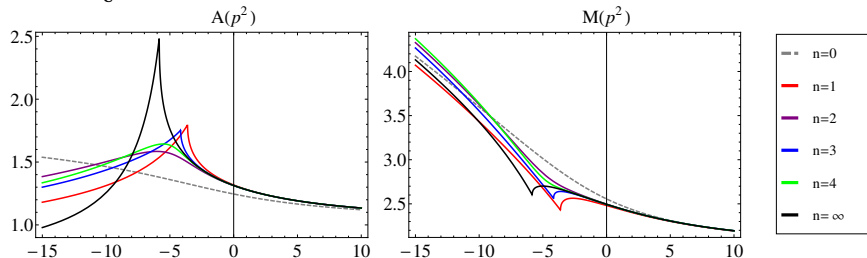


Quark propagator: $S^{-1}(p) = ipA(p^2) + B(p^2)$; $M(p^2) := B(p^2)/A(p^2)$

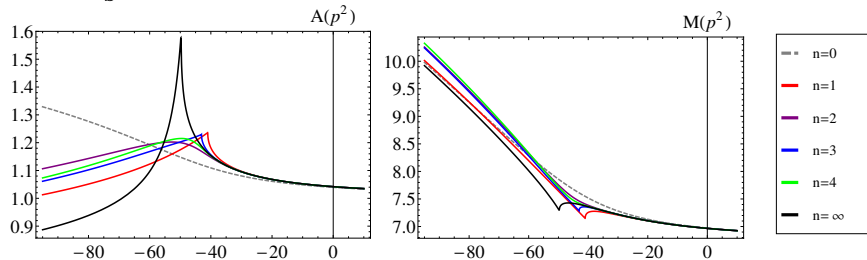
Solutions to the DSE for n loops

for $m_c = 1.33$ GeV:

$C = 0.51$



for $m_b = 4.62$ GeV:



Quark propagator: $S^{-1}(p) = i\cancel{p}A(p^2) + B(p^2)$; $M(p^2) := B(p^2)/A(p^2)$

Dyson-Schwinger/Bethe-Salpeter approach

The interaction model

BSE:

$$\Gamma_M(p; P) = \int \frac{d^4 q}{(2\pi)^4} g^2 [\mathcal{K}(p; q; P)] \underbrace{S(q_+) \Gamma_M(q; P) S(q_-)}_{\chi_M(q; P)} \Gamma_\nu^a(q; p)$$

$$p_+ = \eta P, \quad p_- = (\eta - 1)P$$

Use an appropriate kernel

that guarantees validity of Ward-Takahashi identities

Bender et al. PRC**65**(2002)

Generalize to $m_q \neq m_{\bar{q}}$ and use our MN model to simplify the calculation

Dyson-Schwinger/Bethe-Salpeter approach

The interaction model

Generalize to $m_q \neq m_{\bar{q}}$

$$\Gamma_M(p; P) = \int_q^\Lambda \mathcal{D}_{\mu\nu}(p-q) l^a \frac{1}{2} [\gamma_\mu \chi_M(q; P) l^a \Gamma_\nu(q_-, p_-) + \Lambda_{M\nu}^a(q, p; P) S(q_-) \gamma_\mu + l^a \Gamma_\nu(q_+, p_+) \chi_M(q; P) \gamma_\mu + \gamma_\mu S(q_+) \Lambda_{M\nu}^a(q, p; P)],$$

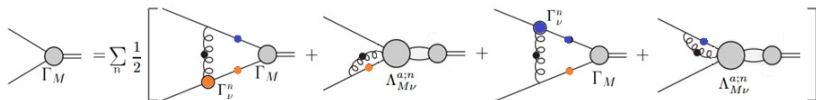


figure adapted from Bender et al. PRC **65** (2002)

Derive recursion relation for $\Lambda_{M\nu}^a(q, p; P)$ for $m_q \neq m_{\bar{q}}$

\Rightarrow Corrections are more involve than for $m_q = m_{\bar{q}}$

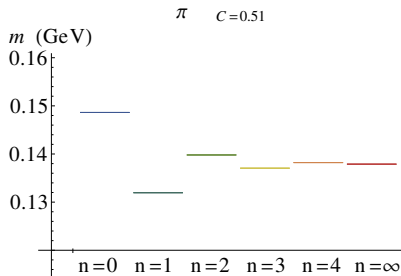
Search for bound-state solutions...

- Pseudoscalar mesons: $q\bar{q}$, $Q\bar{Q}$, $q\bar{Q}$
- Vector mesons: $q\bar{q}$, $Q\bar{Q}$, $q\bar{Q}$
- Scalars \rightarrow no solution
- Axial vectors \rightarrow no solution

Solutions to the BSE

First impression for $q\bar{q}$ systems

The pion



$n = \#$ of loops

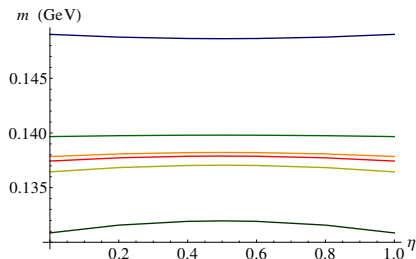
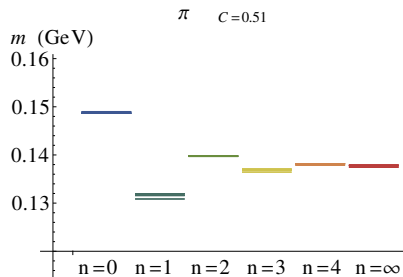
Now: check η -dependence, since this is an issue for unequal-mass constituents

(see also original MN-paper, η was used as fitting parameter).

$$\eta = \frac{p_+}{P}, \quad (\eta - 1) = \frac{p_-}{P}$$

Why η -dependence?

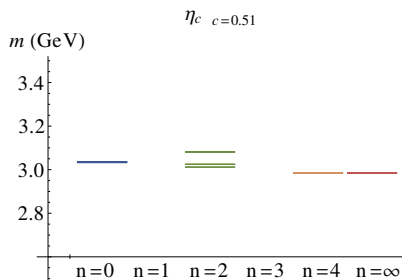
Because MN omits 2 of 4 covariants in BS amplitude \rightarrow this breaks Lorentz covariance: systematic error $< 0.8\%$



Correction to RL $\sim 8\%$ for $\eta = 0.5$

Heavy quarkonia $Q\bar{Q}$

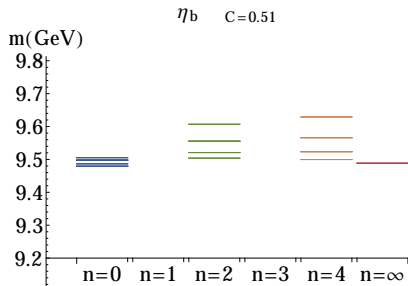
Charmonium



Correction to RL $\sim 1.6\%$ for $\eta = 0.5$
Error due to η -dep $< 6\%$

η -dep increases with increasing n

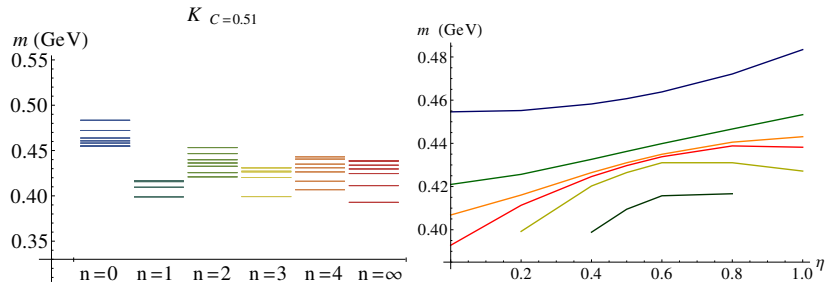
Bottomonium



Correction to RL $\sim 0.2\%$ for $\eta = 0.5$
Error due to η -dep $< 1.5\%$

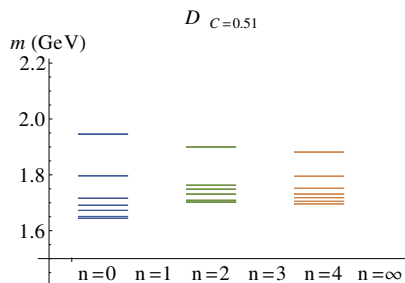
First check for mesons with $m_q \neq m_{\bar{q}}$:

Kaon

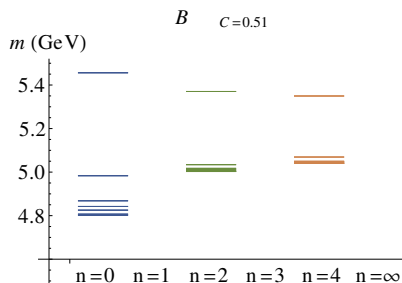


Correction to RL $\sim 7.2\%$ for $\eta = 0.5$
Error due to η -dep $< 10\%$

$q\bar{Q}$ -mesons



Correction to RL $\sim 6.3\%$ for $\eta = 0.75$
Error due to η -dep $< 15\%$
 $m_D^{\text{full}} = 1.68$ GeV



Correction to RL $\sim 3.6\%$ for $\eta = 0.95$
Error due to η -dep $< 12\%$
 $m_B^{\text{full}} = 5.08$ GeV

η -dep decreases with increasing n

Summary

- Interest in $q\bar{Q}$ meson properties, phenomenology and structure from QCD
- Showed you an approach to these problems within DSE/BSE formalism
- Based on previous works, used an interaction model where
 - ★ Very heavy quarks and mesons can be studied numerically (even heavy-quark limit)
 - ★ Effects beyond the popular RL truncation of DSE/BSE can be studied systematically and quantitatively

$$\pi \sim 8\%$$

$$\eta_c \sim 1.6\%$$

$$\eta_b \sim 0.2\%$$

$$K \sim 7.2\%$$

$$D \sim 6.3\%$$

$$B \sim 3.6\%$$

BUT: η -dependence must be better understood

Conclusions and Outlook

- Setup established previously was adapted to investigate $q\bar{q}$, $Q\bar{Q}$, $q\bar{Q}$ meson
- Corrections to the RL truncation are more complicated in $q\bar{Q}$ mesons than in $q\bar{q}$ mesons
- Model artifacts have to be taken into account and kept under control (η -dependence)
- Heavy-light meson masses computed from the given model and parameters are reasonable
- Check heavy-quark symmetry predictions: relations between vector and pseudoscalar mesons in extreme case $m_Q \gg m_q$
- Decay constants

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Thank you!