

# Chiral symmetry, meson spectroscopy, and medium modifications

Thomas Hilger

Karl-Franzens University Graz

Leibnitz

October 7, 2015

together with:

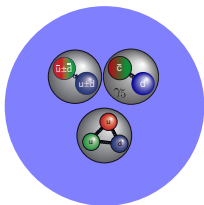
- Maria Gómez-Rocha (ECT\*, Trento), Carina Popovici, Andreas Krassnigg (Univ. Graz), Wolfgang Lucha (HEPHY, Vienna)
- Sergey Dorkin (JINR, Dubna & Univ. Dubna), Leonid Kaptari (JINR, Dubna & HZDR, Dresden), Burkhard Kämpfer (HZDR, Dresden & TU Dresden)
- Stefan Leupold (Univ. Uppsala)

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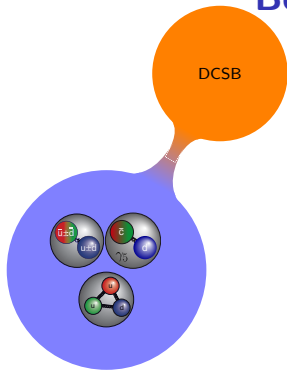
Austrian Science Fund (FWF) project no. P25121-N27,  
Heisenberg-Landau program of the JINR-FRG collaboration,  
GSI-FE, BMBF

[Covariant.ModelsOfHadrons.com](http://Covariant.ModelsOfHadrons.com)

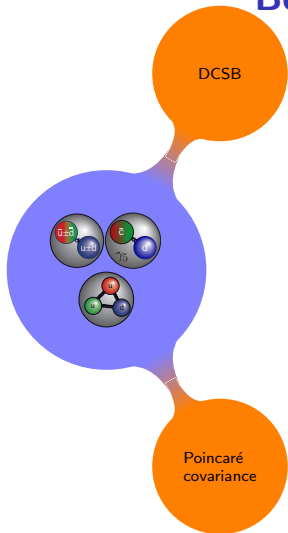
# Meson Spectroscopy and the Dyson-Schwinger– Bethe-Salpeter Equation Approach



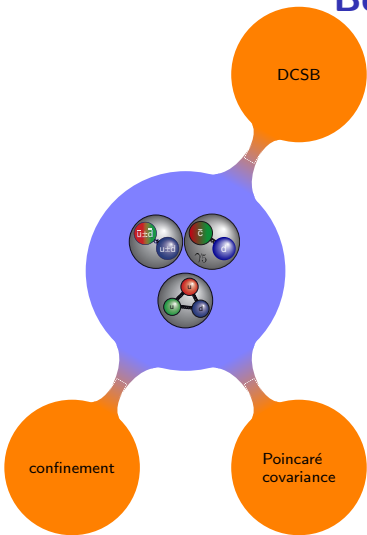
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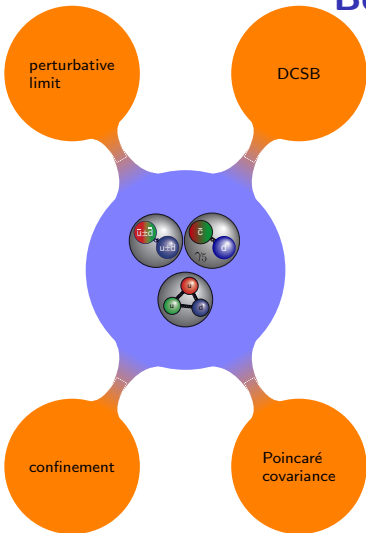
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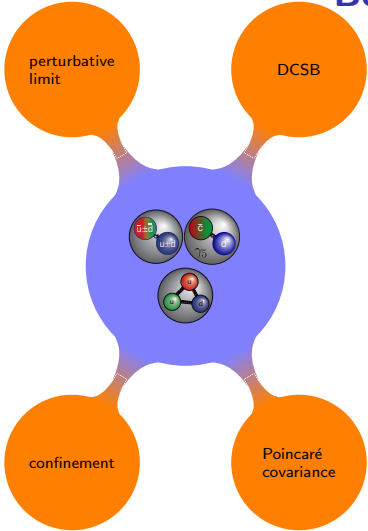
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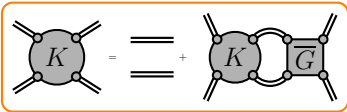
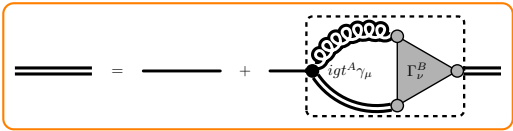
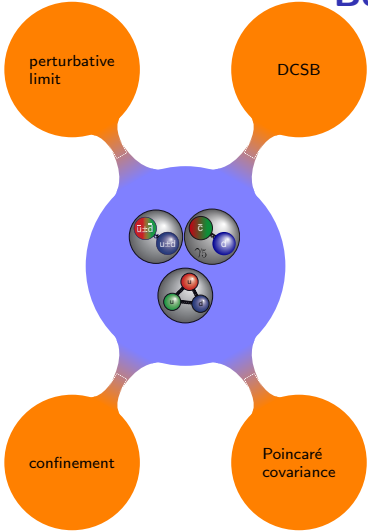
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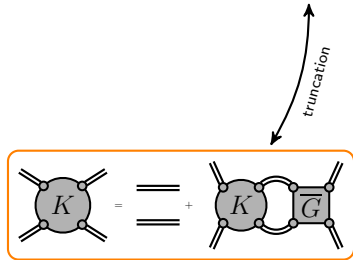
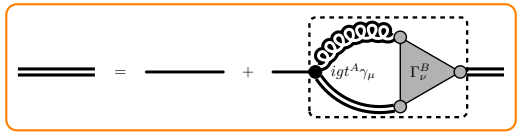
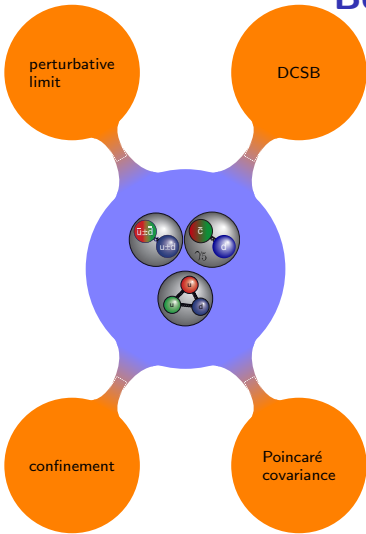
Feynman diagram equation showing a double line (representing a meson) equal to a single line (representing a quark) plus a loop diagram. The loop diagram consists of a wavy line (representing a gluon) and a vertex labeled  $\Gamma_{\nu}^B$ .



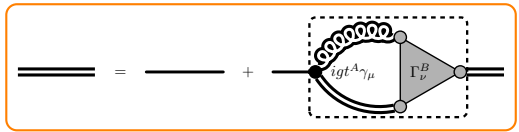
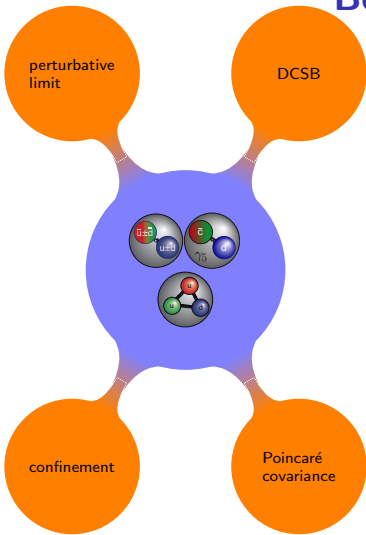
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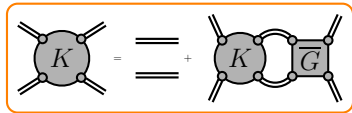


# Meson Spectroscopy and the Dyson-Schwinger–Bethe-Salpeter Equation Approach

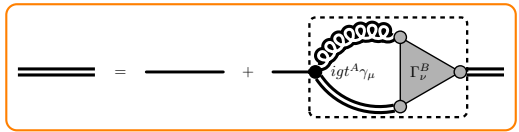
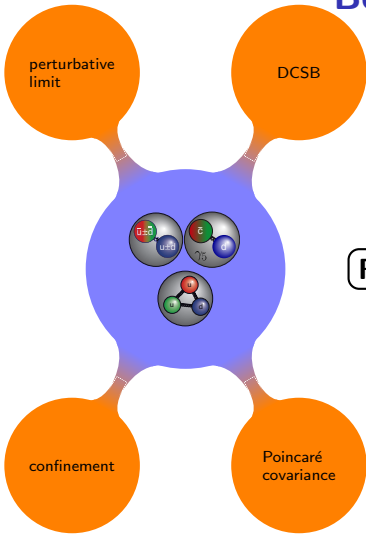


avWTI

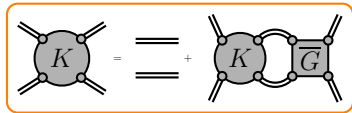
truncation



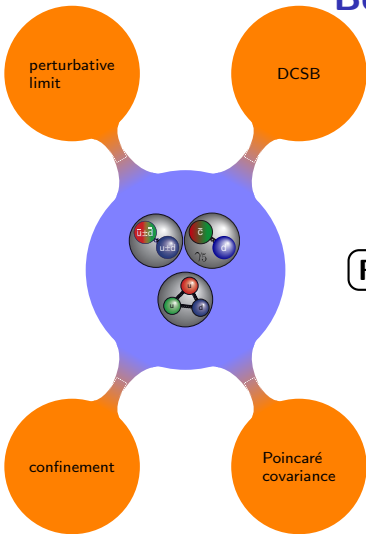
# Meson Spectroscopy and the Dyson-Schwinger–Bethe-Salpeter Equation Approach



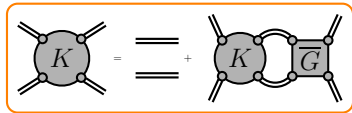
Rainbow-Ladder truncation  $\Rightarrow$  avWTI



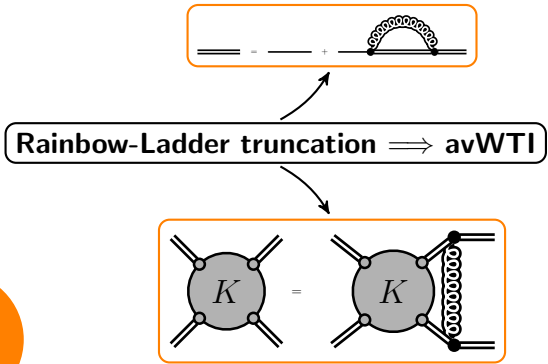
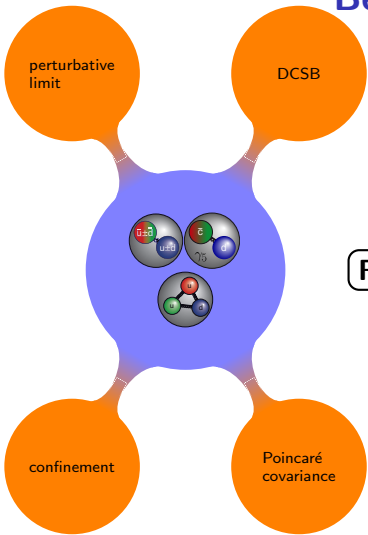
# Meson Spectroscopy and the Dyson-Schwinger–Bethe-Salpeter Equation Approach



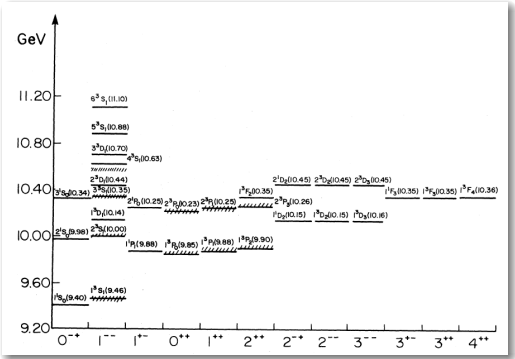
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# Meson Spectroscopy and the Dyson-Schwinger–Bethe-Salpeter Equation Approach



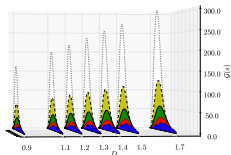
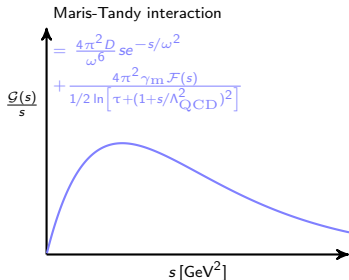
# Incitement



[Bottomonium by Godfrey, Isgur, 1985]

# Model and Strategy

... so far no comprehensive attempt at RL meson phenomenology



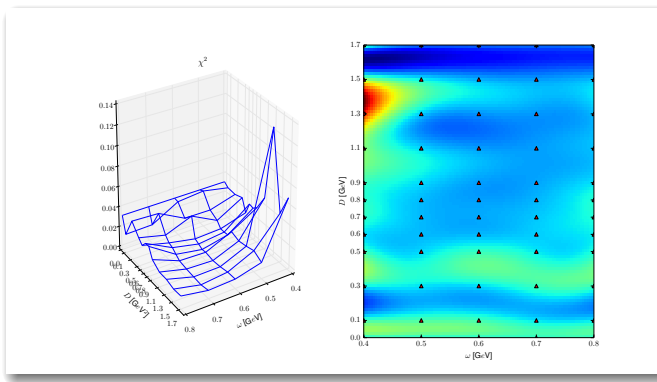
... is that good enough?

- application to systems where corrections to RL are expected to be least important → bottomonium
- leave functional and UV form unchanged
- allow for more freedom in the effective interaction → quark mass dependence, vary  $\omega$  and  $D$  independently
- include lowest radial excitations
- $J = 0, 1, 2, \dots$



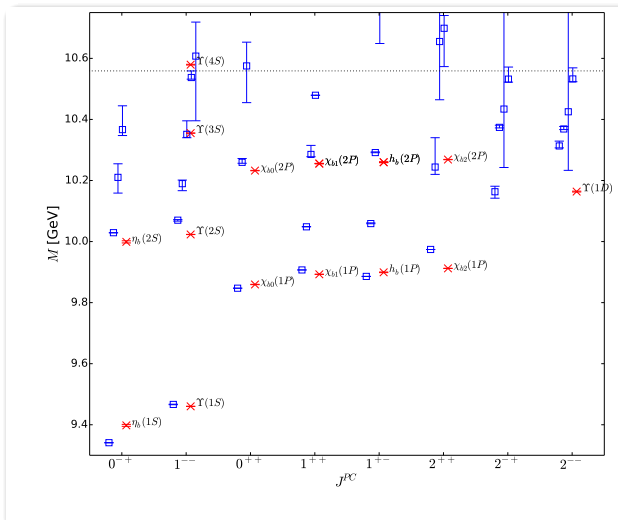
# Bottomonium

- evaluate splittings at  $(\omega - D)$ -grid
- find minimal  $\chi^2(\omega, D) = \sum_{\text{splittings}} (\Delta M_{\text{exp}} - \Delta M_{\text{th}})^2$
- find minimal  $\bar{\chi}^2(m_q) = \sum_{\text{groundstates}} (M_{\text{exp}} - M_{\text{th}})^2$  for optimal  $(\omega, D)$



[C. Popovici, T. Hilger, M. Gómez-Rocha, A. Krassnigg, Few Body Syst. **56**: 481, 2015.]

# Bottomonium



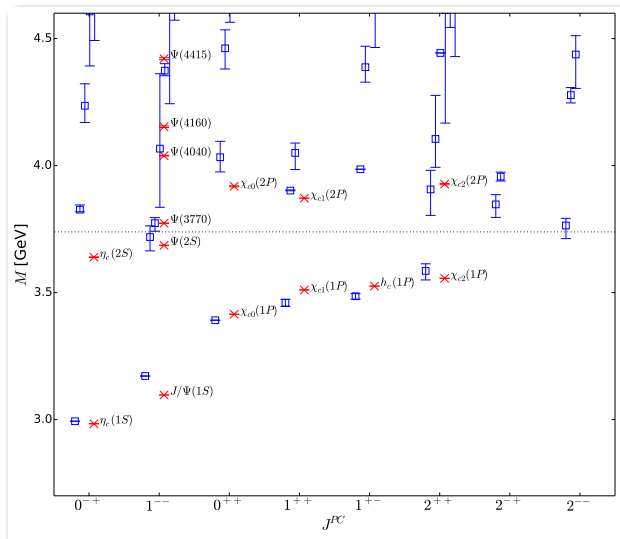
× experiment

- good identification of states
- well reproduced splittings (excitations, level orderings)

[T. Hilger, C. Popovici, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 034013, 2015.]

●  $m_b = 3.635$  GeV at  $\mu = 19$  GeV,  $\omega = 0.7$  GeV,  $D = 1.3$  GeV<sup>2</sup>

# Charmonium



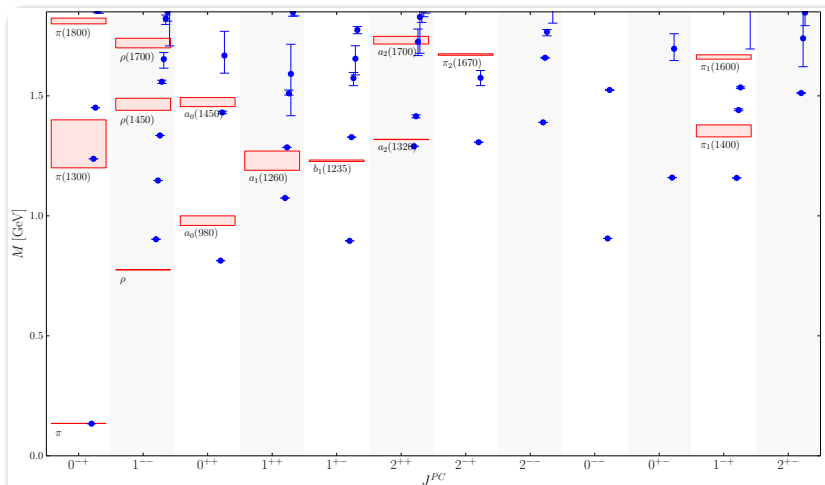
× experiment

- no extra states
- excellently reproduced splittings, in particular  $1^{--}$

[T. Hilger, C. Popovici, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 034013, 2015.]

●  $m_c = 0.855$  GeV at  $\mu = 19$  GeV,  $\omega = 0.7$  GeV,  $D = 0.5$  GeV<sup>2</sup>

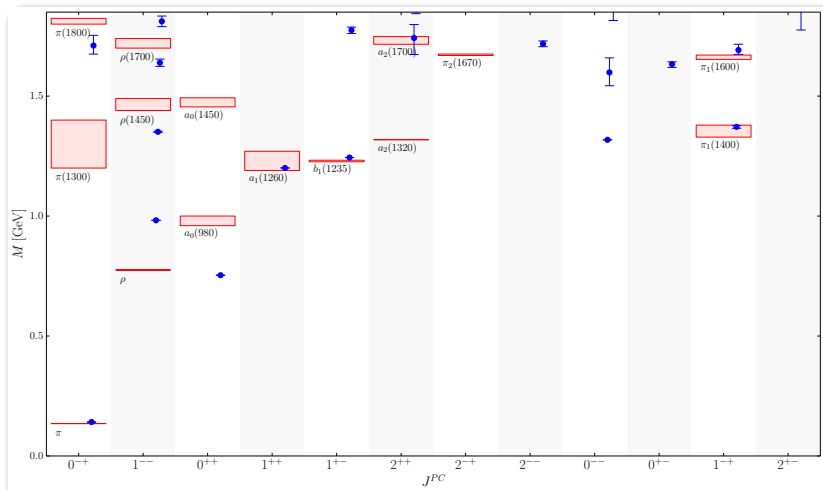
# Light Isovector Quarkonium



[T. Hilger, M. Gómez-Rocha, A. Krassnigg, arXiv:1508.07183]

●  $m_q = 0.003$  GeV at  $\mu = 19$  GeV,  $\omega = 0.4$  GeV,  $D = 1.7$  GeV<sup>2</sup>

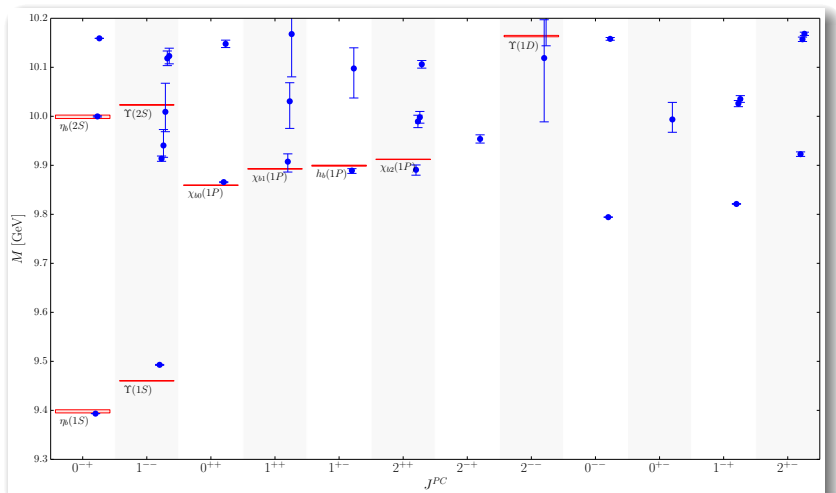
# Exotics: Light Isovector Quarkonium



[T. Hilger, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 114004, 2015.]

●  $m_q = 0.003$  GeV at  $\mu = 19$  GeV,  $\omega = 0.7$  GeV,  $D = 1.4$  GeV<sup>2</sup>

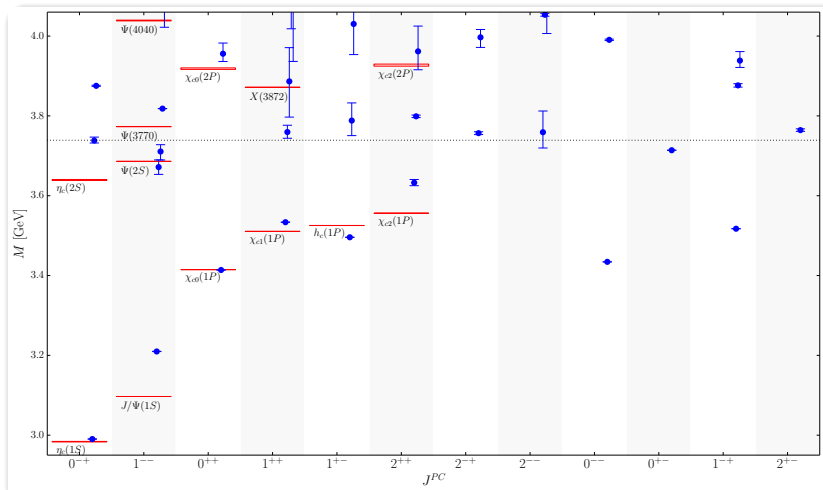
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●  $m_b = 3.635$  GeV at  $\mu = 19$  GeV,  $\omega = 0.7$  GeV,  $D = 0.8$  GeV<sup>2</sup>

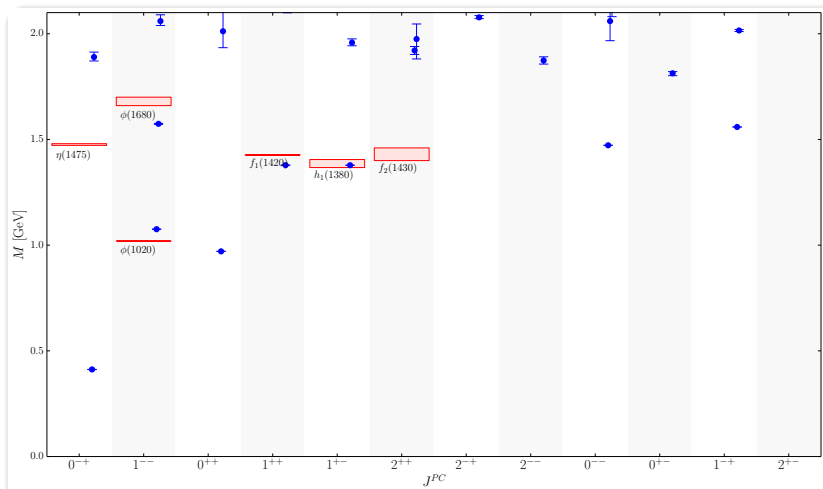
# Exotics: Charmonium



[T. Hilger, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 114004, 2015.]

●  $m_c = 0.855$  GeV at  $\mu = 19$  GeV,  $\omega = 0.6$  GeV,  $D = 0.9$  GeV<sup>2</sup>

# Exotics: Strangeonium

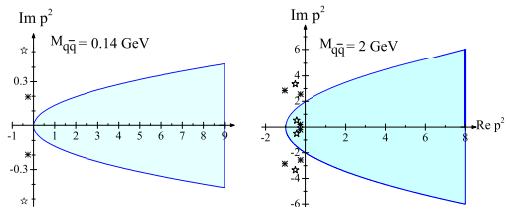


[T. Hilger, M. Gómez-Rocha, A. Krassnigg, arXiv:1508.07183]

●  $m_q = 0.070$  GeV at  $\mu = 19$  GeV,  $\omega = 0.8$  GeV,  $D = 1.7$  GeV<sup>2</sup>



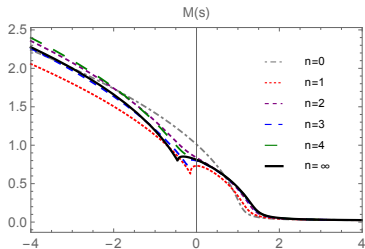
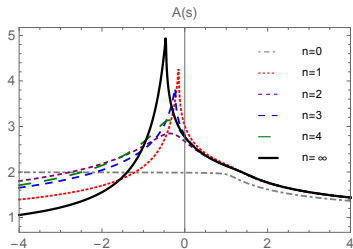
# A Path to Open Flavor Meson Spectroscopy



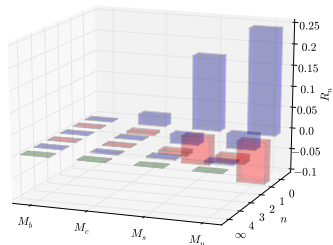
[S. Dorkin, L. Kaptari, T. Hilger, B. Kämpfer, Phys. Rev. C **89**: 034005, 2014.]

- numerics technically involved
- Outlook: AVWTI construction
- estimate BRL corrections
- Munczek-Nemirovsky model:  $\mathcal{G}(s) \propto s\delta^{(4)}(s)$
- integral equations reduce to algebraic equations
- infinite dressing of quark-gluon vertex with gluon loops possible

# Quark-Gluon Vertex and Quark Propagator



[M. Gómez-Rocha, T. Hilger, A. Krassnigg, accepted by PRD, arXiv:1506.03686] [M. Gómez-Rocha, T. Hilger, A. Krassnigg, Few Body Syst. **56**: 475, 2015.]

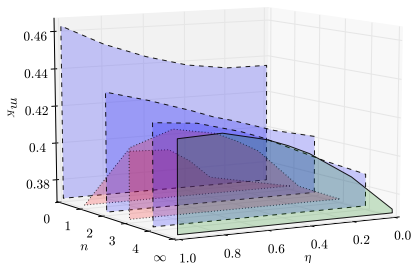


$$R_n := \frac{M_n(s=0) - M_\infty(s=0)}{M_\infty(s=0)}$$

- qualitative differences on timelike domain
- weaker effect for heavier quarks

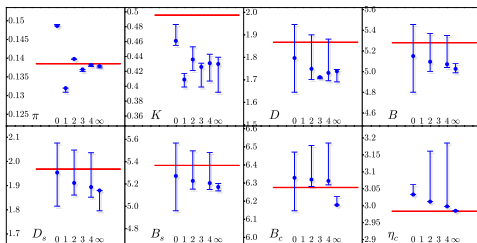
# Quark-Gluon Vertex Dressing and Meson Masses

[M. Gómez-Rocha, T. Hilger, A. Krassnigg, accepted by PRD, arXiv:1506.03686]



- momentum partitioning dependence
- minimize dressing correction
- error estimate

- sizeable but not overwhelming dressing effects
- careful, comprehensive RL phenomenology worthwhile



# Reminder: Chiral Symmetry

two flavor Lagrangian:

$$\mathcal{L} = \begin{pmatrix} \bar{u} \\ \bar{d} \end{pmatrix}^T \left( i\gamma_\mu \partial^\mu - \begin{bmatrix} m_u & 0 \\ 0 & m_d \end{bmatrix} \right) \begin{pmatrix} u \\ d \end{pmatrix}$$

$P = -1$  meson current:

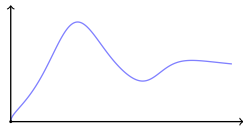
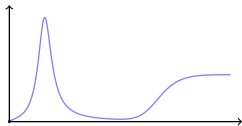
$$j_\mu^{V,\tau}(x) = \bar{\psi} \gamma_\mu \tau \psi$$

$P = +1$  meson current:

$$j_\mu^{A,\tau}(x) = \bar{\psi} \gamma_5 \gamma_\mu \tau \psi$$

current-current correlator:

$$\Pi_{\mu\nu}^X(q) = i \int d^4x e^{-iqx} \langle 0 | T [j_\mu^{X,\tau}(x) (j_\nu^{X,\tau}(0))^\dagger] | 0 \rangle$$



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chiral transformations:

$$\mathcal{T}: \psi_{R,L} \equiv \frac{1 \pm \gamma_5}{2} \begin{pmatrix} u \\ d \end{pmatrix} \rightarrow e^{-i\frac{\vec{\lambda}}{2} \vec{\Theta}_{R,L}} \psi_{R,L}$$

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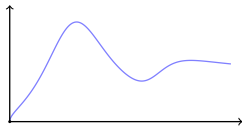
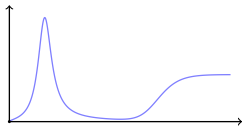
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invariant for  $m_{u,d} = 0$

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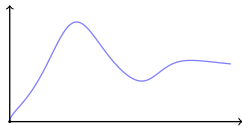
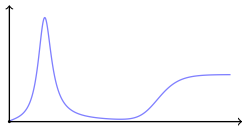
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# Reminder: Chiral Symmetry

two flavor Lagrangian:

$$\mathcal{L} = \left(\frac{\bar{u}}{d}\right)^T \left( i\gamma_\mu \partial^\mu - \begin{bmatrix} m_u & 0 \\ 0 & m_d \end{bmatrix} \right) \begin{pmatrix} u \\ d \end{pmatrix}$$

chiral transformations:

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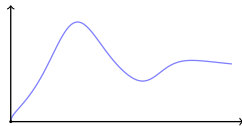
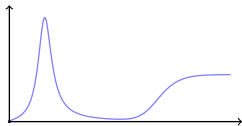
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conversion by set of finite  $\{\Theta_{R,L}\}$

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 $j_\mu^{A,\tau}(x) = \bar{\psi} \gamma_5 \gamma_\mu \tau \psi$

current-current correlator:

$$\Pi_{\mu\nu}^X(q) = i \int d^4x e^{-iqx} \langle 0 | T [j_\mu^{X,\tau}(x) (j_\nu^{X,\tau}(0))^\dagger] | 0 \rangle$$



# Reminder: Chiral Symmetry

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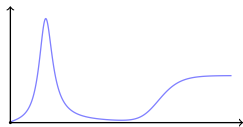
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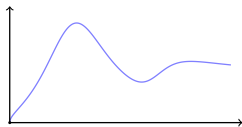
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chirally symmetric  $|0\rangle$

parity blind correlators  
degenerate spectra





# Reminder: Chiral Symmetry and Open Flavor Mesons

three flavor Lagrangian:

$$\mathcal{L} = \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{h} \end{pmatrix}^T \left( i\gamma_\mu \partial^\mu - \begin{bmatrix} m_u & 0 & 0 \\ 0 & m_d & 0 \\ 0 & 0 & m_h \end{bmatrix} \right) \begin{pmatrix} u \\ d \\ h \end{pmatrix}$$

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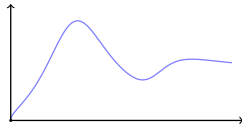
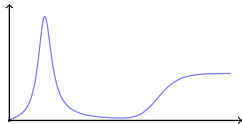
$$j_{\mu}^{V, \tau}(x) = \bar{\psi} \gamma_{\mu} \tau \psi$$

$P = +1$  meson current:

$$j_{\mu}^{A, \tau}(x) = \bar{\psi} \gamma_5 \gamma_{\mu} \tau \psi$$

current-current correlator:

$$\Pi_{\mu\nu}^X(q) = i \int d^4x e^{-iqx} \langle 0 | T [j_{\mu}^{X, \tau}(x) (j_{\nu}^{X, \tau}(0))^{\dagger}] | 0 \rangle$$



# Reminder: Chiral Symmetry and Open Flavor Mesons

three flavor Lagrangian:

$$\mathcal{L} = \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{h} \end{pmatrix}^T \left( i\gamma_\mu \partial^\mu - \begin{bmatrix} m_u & 0 & 0 \\ 0 & m_d & 0 \\ 0 & 0 & m_h \end{bmatrix} \right) \begin{pmatrix} u \\ d \\ h \end{pmatrix}$$

chiral transformations:

$$\mathcal{T} : \psi_{R,L} \equiv \frac{1 \pm \gamma_5}{2} \begin{pmatrix} u \\ d \\ h \end{pmatrix}$$

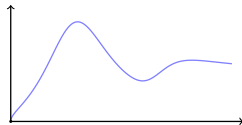
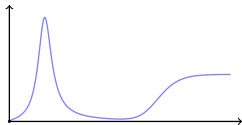
$$\rightarrow \exp \left\{ -\frac{i}{2} \begin{bmatrix} \Theta_3 & \Theta_1 - i\Theta_2 & 0 \\ \Theta_1 + i\Theta_2 & -\Theta_3 & 0 \\ 0 & 0 & 0 \end{bmatrix}_{R,L} \right\} \psi_{R,L}$$

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# Reminder: Chiral Symmetry and Open Flavor Mesons

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invariant for  $m_{u,d} = 0$

chiral transformations:

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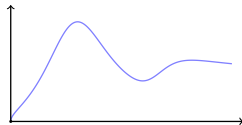
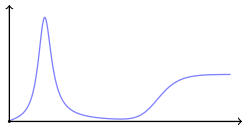
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# Reminder: Chiral Symmetry and Open Flavor Mesons

three flavor Lagrangian:

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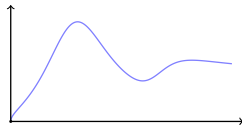
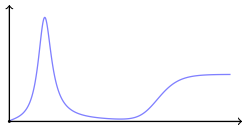
conversion by set of finite  $\{\Theta_{R,L}\}$

$P = -1$  meson current:  
 $j_{\mu}^{V,\tau}(x) = \bar{\psi} \gamma_\mu \tau \psi$

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current-current correlator:

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# Reminder: Chiral Symmetry and Open Flavor Mesons

three flavor Lagrangian:

$$\mathcal{L} = \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{h} \end{pmatrix}^T \left( i\gamma_\mu \partial^\mu - \begin{bmatrix} m_u & 0 & 0 \\ 0 & m_d & 0 \\ 0 & 0 & m_h \end{bmatrix} \right) \begin{pmatrix} u \\ d \\ h \end{pmatrix}$$

invariant for  $m_{u,d} = 0$

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$$\rightarrow \exp \left\{ -\frac{i}{2} \begin{bmatrix} \Theta_3 & \Theta_1 - i\Theta_2 & 0 \\ \Theta_1 + i\Theta_2 & -\Theta_3 & 0 \\ 0 & 0 & 0 \end{bmatrix}_{R,L} \right\} \psi_{R,L}$$

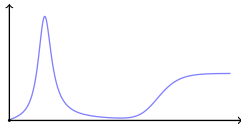
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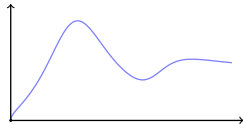
current-current correlator:

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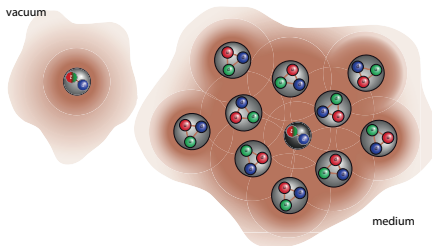
chirally symmetric  $|0\rangle$

parity blind correlators degenerate spectra



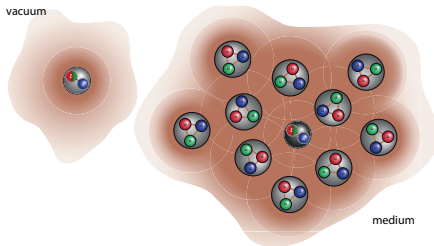
# Open Flavor Mesons in the Medium

- A Window to DCSB



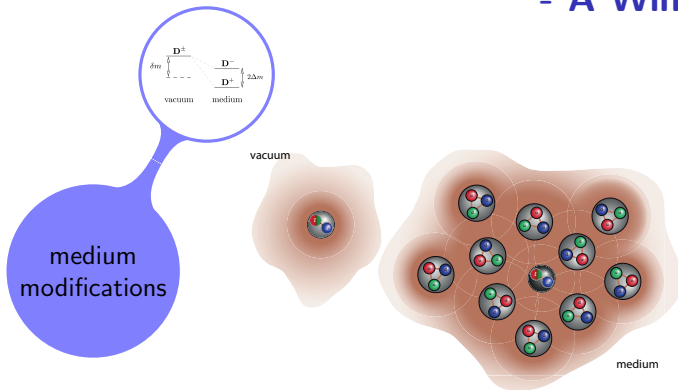
# Open Flavor Mesons in the Medium

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# Open Flavor Mesons in the Medium

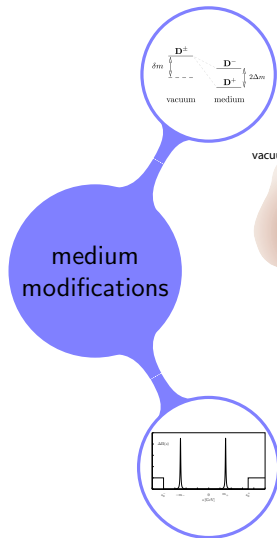
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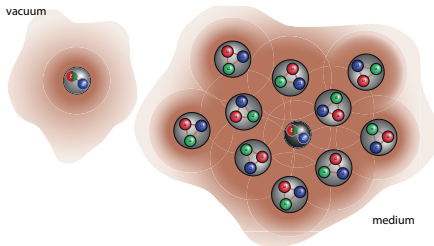


# Open Flavor Mesons in the Medium

- A Window to DCSB

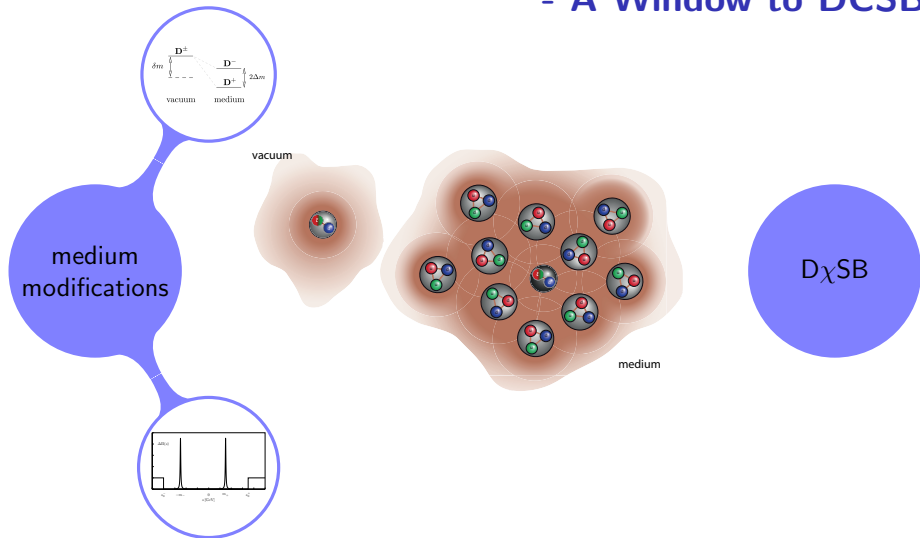


medium  
modifications



# Open Flavor Mesons in the Medium

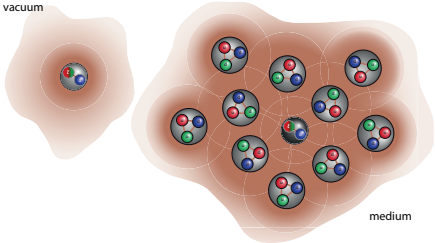
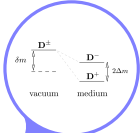
- A Window to DCSB



# Open Flavor Mesons in the Medium

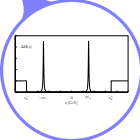
## - A Window to DCSB

medium modifications



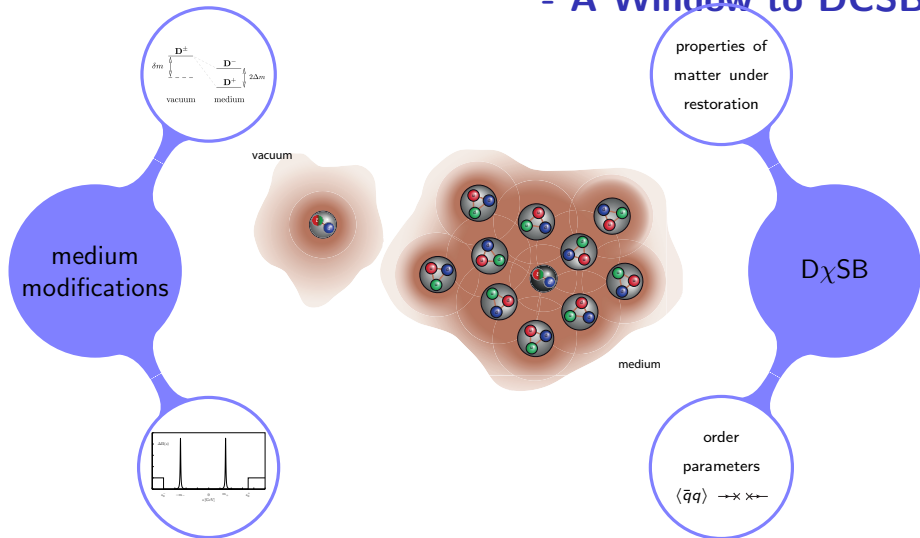
properties of matter under restoration

$D\chi SB$



# Open Flavor Mesons in the Medium

## - A Window to DCSB



# Hadron physics and QCD sum rules

current-current correlator

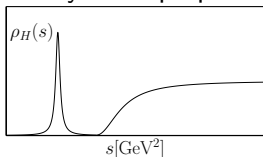
$$\Pi_{\mu\nu}(q) = i \int d^4x e^{iqx} \langle T [j_\mu(x) (j_\nu(0))^\dagger] \rangle$$

dispersion relation

$$\Pi(q^2) = \frac{1}{\pi} \int_0^\infty ds \frac{\Delta\Pi(s)}{s-q^2}$$



spectral density  $\leftrightarrow$  hadronic properties



separation of scales

operator product expansion

$$= C_1(q) + C_2(q) \langle \bar{q}q \rangle + C_3(q) \langle \bar{q}g\sigma\mathcal{G}q \rangle + \dots$$

$$= \text{[circle with arrow]} + \text{[circle with arrow and cross]} + \text{[circle with arrow and cross and vertical line with cross]} + \dots$$

QCD condensates:

- encode medium dependence
- order parameters of chiral symmetry phase transition

$$\int_0^\infty ds \rho_H(s) = \text{[circle with arrow]} + \text{[circle with arrow and cross]} + \text{[circle with arrow and cross and vertical line with cross]} + \dots$$

# Order Parameters and Light Quark Currents

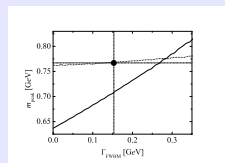
- $\langle \bar{q}q \rangle$  suppression in light-quark meson operator product expansion (e.g.  $\rho$  meson sum rules):  $m_q \langle \bar{q}q \rangle$
- $\langle \bar{q}q \rangle$  influence via assumptions/models: e.g.
  - $\langle \bar{q}\Gamma q \bar{q}\Gamma q \rangle \propto \langle \bar{q}q \rangle^2$   
→ fragile transition to medium
  - continuum threshold  $s_0 \leftrightarrow f_\pi \leftrightarrow \langle \bar{q}q \rangle$
- determination of other order parameters (e.g. four-quark condensates  $\langle \bar{q}\Gamma q \bar{q}\Gamma q \rangle$ ) is model dependent

# Properties of the $\rho$ Meson under Chiral Symmetry Restoration

impact of  $D\chi$ SB order parameters on the  $\rho$  meson and implications of chirally symmetric sum rules

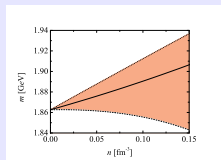
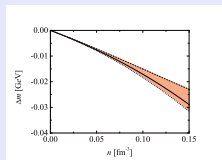
[TH, R. Thomas, B. Kämpfer, S. Leupold, Phys. Lett. B **709** (2012)

200]



# Spin-0 Open Flavor Mesons at Finite Density

- particle-antiparticle **splitting** + **shift**



- $\langle \bar{q}q \rangle$  amplification due to heavy charm quark mass:  $m_c \langle \bar{q}q \rangle$

[TH, R. Thomas, B. Kämpfer, Phys. Rev. C **79** (2009) 025202]

[S. Zschocke, TH, B. Kämpfer, Eur. Phys. J. A **47** (2011) 151]

[TH, B. Kämpfer, Nucl. Phys. B Proc. Suppl. **207-208** (2010) 025202]

[TH, B. Kämpfer, Conf. Proc. Italian Phys. Soc. **99** (2010)]

[B. Kämpfer, TH, H. Schade, R. Schulze, G. Wolf, PoSBormio **2010**]

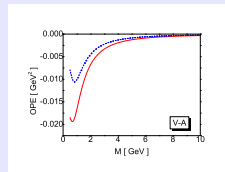
[R. Rapp et al., *In-medium excitations*, Lect. Notes Phys. **814** 335 (2011)]



# Chiral QCD Sum Rules for Open Flavor Spin-0 and -1 Mesons at Finite Density

- Weinberg-Kapusta-Shuryak sum rules:  $\langle \bar{q}q \rangle$ -suppression by light quark mass
- open flavor chiral partner sum rules: spectral differences driven by order parameters only
- amplification of order parameters due to heavy charm quark mass
- hierarchy of order parameters

[TH, B. Kämpfer, S. Leupold, Phys. Rev. C **84** (2011) 045202]  
[TH, T. Buchheim, B. Kämpfer, S. Leupold, Prog. Part. Nucl. Phys. **67** (2012) 188]  
[TH, R. Schulze, B. Kämpfer, J. Phys. G: Nucl. Part. Phys. **37** (2010) 094054]  
[TH, B. Kämpfer, Nucl. Phys. Proc. Suppl. **207-208** (2010) 277]



# Four-Quark Condensates and Open Flavor Mesons

[T. Buchheim, TH, B. Kämpfer,

Phys. Rev. C **91** (2015)] [T. Buchheim,

TH, B. Kämpfer, arXiv:1509.06144]

- order parameters
- spin-0 and -1
- in-medium
- chiral sum rules

# Four-Quark Condensates and Open Flavor Mesons

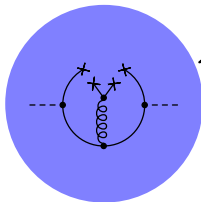
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Phys. Rev. C **91** (2015)] [T. Buchheim,

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$$\langle \bar{\psi} \sigma \tau \Gamma \psi \bar{\psi} \sigma' \tau' \Gamma' \psi \rangle$$



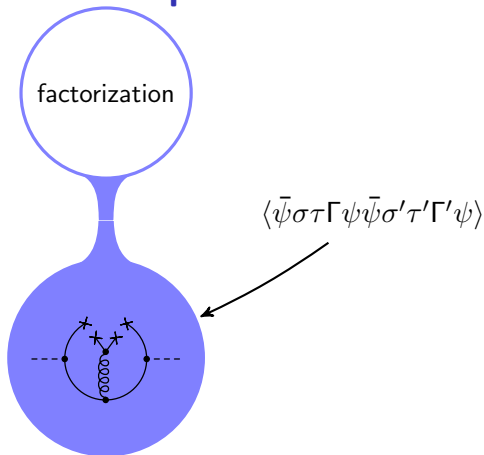
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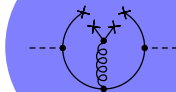
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factorization

$$\langle \bar{\psi} \sigma \tau \Gamma \psi \bar{\psi} \sigma' \tau' \Gamma' \psi \rangle$$



heavy-  
quark mass  
expansion

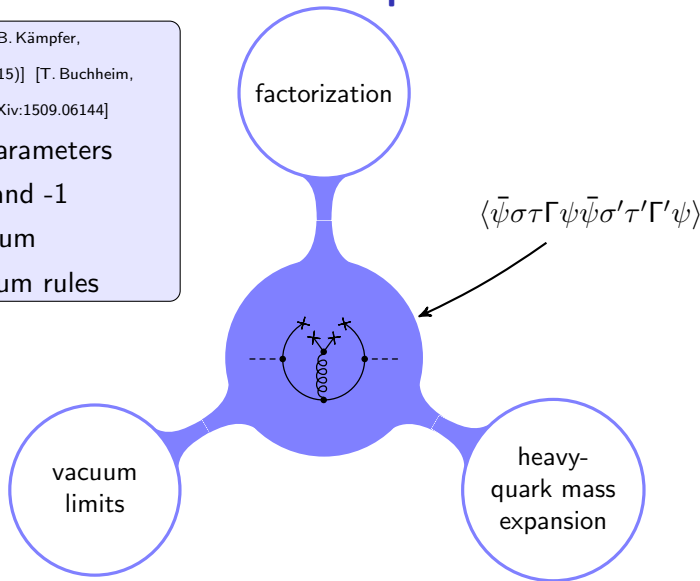
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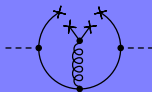
Phys. Rev. C **91** (2015)] [T. Buchheim,

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vacuum  
limits

heavy-  
quark mass  
expansion

[T. Buchheim, TH, B. Kämpfer, J. Phys. Conf. Ser. **503** (2014)]

[T. Buchheim, TH, B. Kämpfer, E.P.J. WoC **81** (2014)]

[T. Buchheim, TH, B. Kämpfer, Nucl. Phys. Proc. Suppl. **258-259** (2015)]

# Summary I

- quark mass dependence of effective interaction
- optimized rainbow-ladder DS-BS study describes ground states and lowest radial excitations
- extra states in vector- and axial-vector channel for bottomonium
- *exotic* charmonium, bottomonium (and light isovector) spectrum
- 1st BRL study of open flavor mesons in Munczek-Nemirovsky model
- stay tuned: ... decay constants ... open flavor mesons ...

[M. Gómez-Rocha, T. Hilger, A. Krassnigg, accepted by PRD, arXiv:1506.03686]

[T. Hilger, M. Gómez-Rocha, A. Krassnigg, arXiv:1508.07183]

[T. Hilger, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 114004, 2015.]

[T. Hilger, C. Popovici, M. Gómez-Rocha, A. Krassnigg, Phys. Rev. D **91**: 034013, 2015.]

[M. Gómez-Rocha, T. Hilger, A. Krassnigg, Few Body Syst. **56**: 475, 2015.]

[C. Popovici, T. Hilger, M. Gómez-Rocha, A. Krassnigg, Few Body Syst. **56**: 481, 2015.]



# Summary II

- light four-quark condensates for  $D$  mesons in medium

[T. Buchheim, TH, B. Kämpfer, Phys. Rev. C **91** (2015)]

[T. Buchheim, TH, B. Kämpfer, arXiv:1509.06144]

- heavy-quark mass expansion and factorization of four-quark condensates in the medium
- continuous transition from medium to vacuum  $\rightarrow$  algebraic vacuum limits

[T. Buchheim, TH, B. Kämpfer, J. Phys. Conf. Ser. **503** (2014)]

[T. Buchheim, TH, B. Kämpfer, Nucl. Phys. Proc. Suppl. **258-259** (2015)]

[T. Buchheim, TH, B. Kämpfer, E. P. J. WoC **81** (2014)]

Thank You!

Your funding  
agency  
could be here!